# SOME GEOMETRIC ASPECTS OF GRAPHS AND THEIR EIGENFUNCTIONS

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#### Abstract

We study two mathematical notions, that of nodal regions for eigenfunctions of the Laplacian, and that of fiber products, in the context of graph theory. We formulate analogous notions and theorems for graphs and their eigenpairs. These techniques suggest new ways of studying problems related to spectral theory of graphs.

#### 1 Introduction

In analysis on manifolds there is an extensive literature on isoperimetric problems and the Laplacian. Some analogues in graph theory, usually concerning eigenvalues of the adjacency matrix or the associated "Laplacian," are known (see [Alo86, CDS79, Dod84]). But on the whole much less in known for graphs, especially for isoperimetric type problems, and many tools from analysis are in want of a good generalization to graph theory.

In this paper we show that the concept of nodal regions in analysis has a precise analogue in graph theory. This gives us geometric insight into the eigenvectors. We show how this, along with information theory and graph coverings, can give some slight improvements to certain eigenvalue bounds. We also show that the mathematical concept of a fiber product gives an interesting type of graph product; it generates new d-regular graphs from old ones in a simple manner, and numerical experiments show that it can yield graphs with small second eigenvalue.

In section 2 we notice that graph eigenvectors can be viewed as minimizers of a Rayleigh quotient over, say, piecewise differentiable functions on the geometric realization of the graph. This suggests a notion of "graph with boundary" and what their adjacency matrices should

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be. All the standard comparison theorems about eigenvalues of the Laplacian and nodal regions of eigenfunctions of the Laplacian carry over verbatim to graphs.

In particular there is a precise graph analogue of the fact that when Dirichlet eigenfuctions of the Laplacian on manifolds are restricted to any of their nodal regions, they give the first Dirichlet eigenfunction of that region. We use this property to study some aspects of eigenfunctions and eigenvalues in what follows. We show, in section 3, that a d-regular graph for  $d \geq 3$  of diameter 2k has a second eigenvalue of at least  $2\sqrt{d-1}\left(1-\pi^2/(2k^2)+O(k^4)\right)$ ; the proof can be stated in elementary terms, but we use the language of section 2 partly to point out the relationship with the classical eigenvalue estimating techniques.

In section 4 we discuss isoperimetric aspects of subgraphs of the d-regular infinite tree. Here the situation is much different than the analogue, in analysis, namely for subdomains of  $\mathbb{R}^n$ , at least for classical isoperimetric problems. However, we conjecture that an analogue of the Faber-Krahn inequality holds for subgraphs of the d-regular infinite tree, and prove a weaker form of the conjecture. In section 5 we discuss the problem of when a second eigenvalue can "persist" in covers of the graphs. We give an example of a case where the second eigenvalue persists in an infinite set of covers. Such examples yield somewhat unusual nodal regions, and we conjecture that this behavior is exceptional. We remark on its connections to "Ramanujan graphs" and number theory.

In section 6 we make a numerical study of the fiber product operation (explained in the appendix), on some very simple "Ramanujan graphs." The fiber product is quite simple to work with (e.g. on a computer). By forming "twisted" fiber products we make the empirical observation that twisting sometimes reduces the second eigenvalue, yielding graphs with comparitively small second eigenvalue (e.g. it seems hard to find graphs of the same number of vertices and degree with second eigenvalue as small). We compare these to random graphs and graphs obtained by heuristic search, and in the process note some interesting properties of the random graphs' second eigenvalue.

In section 7 we discuss other directions suggested by this work and some questions unsolved in this investigation.

In appendix A we gather some background for the techniques used in this paper. In the first part we review Shannon's algorithm for capacity calculations. In the latter two sections we explain the eigenvalue comparison theorems and Galois/covering theory in the context needed here; for the former topic this is virtually identical to the classical theory (from analysis). For the latter topic we summarize the aspects of Galois/covering theory of interest to spectral theory of graphs; this theory is well-known, but is seldom stressed in the graph theory literature, and in particular the notion of fiber product, which arises there, does not seemed to have received the attention it merits (at least not in spectral theory).

Throughout the paper we use the following convention: given an undirected graph, G, on n nodes, we denote by  $\lambda_i = \lambda_i(G)$  the i-th eigenvalue of the adjacency matrix of G arranged in non-increasing order, i.e.  $\lambda_1 \geq \lambda_2 \geq \cdots \geq \lambda_n$ . By the second eigenvalue, denoted  $\lambda = \lambda(G)$ , we mean the largest in absolute value among  $\lambda_2, \ldots, \lambda_n$ ; the spectral radius, denoted  $\rho = \rho(G)$ , is just  $|\lambda|$ . When we speak of second eigenvalue, etc. it will always be the case that G is a d-regular graph and so  $\lambda_1 = d$ .

It is well know that many theorems about eigenvalues of Laplacians on manifolds have graph analogues. Also certain more geometric notions, such as potential theory, have been traslated and applied to classical graph situations (typically to analyze infinite graphs or trees, such as in [Car72], [Dod84]). Part of what we do here is to point out that one sometimes needs to enlarge the class of graphs under consideration to those with boundaries<sup>1</sup>, and that using this one can view graph eigenvectors as eigenfunctions on a geometric object. Such concepts, and the concept of fiber products, give rise to notions which can shed light on questions concerning the eigenvalues. We hope that more can be learned from this approach, and that this might be helpful in studying eigenvalue and isoperimetric type problems on graphs.

## 2 Nodal Regions of Graph Eigenfunctions

We recall the standard definitions of graph eigenvalues and eigenvectors (see [Alo86] or [CDS79] for details). Let G = (V, E) be an undirected graph. The Laplacian of G is the matrix  $\Delta = D - A$ , where A is the adjacency matrix of G and D is the diagonal matrix whose diagonal entries are the degrees of the vertices; the non-negativity of its eigenvalues,  $0 = \nu_1 \le \nu_2 \le \cdots \le \nu_n$  follows from that of its associated Rayleigh quotient on real-valued functions f on V,

$$\mathcal{R}_{\Delta}(f) \equiv \frac{\left(\Delta f, f\right)}{\left(f, f\right)} = \frac{\sum_{(u, v) \in E} \left(f(u) - f(v)\right)^2}{\sum_{v \in V} \left(f(v)\right)^2} \ge 0,$$

where (,) is the  $L^2(V)$  inner product, the above equality being an easy calculation. If G is d-regular, i.e. each vertex has degree d, then A and  $\Delta$  have the same eigenvectors, and the eigenvalues of A,  $d = \lambda_1 \ge \cdots \ge \lambda_n$  satisfy  $\nu_i = d - \lambda_i$ .

Let G = (V, E) be an undirected graph. The geometric realization of G is the metric space  $\mathcal{G}$  consisting of V along with an interval of length 1 glued in between u and v for every edge  $(u, v) \in E$ . By abuse of notation we speak of edges  $e \in E$  and vertices  $v \in V$  as their images in  $\mathcal{G}$ .  $\mathcal{G}$  looks like a one-dimensional manifold except at the vertices; in particular, we can define the notion of differentiability of a function at points in  $\mathcal{G} - V$ .

Let  $\mu_1$  be the measure on  $\mathcal{G}$  which counts the number of vertices, and let  $\mu_2$  be the measure which restricts to Lebesgue measure on each edge interval. Let H denote the set of continuous functions on  $\mathcal{G}$  which are differentiable on all but a finite subset, S, of points of  $\mathcal{G}$  (where S contains V). For  $f \in H$  we define its Rayleigh quotient to be

$$\mathcal{R}(f) = \frac{\int_{\mathcal{G}} |\nabla f|^2 d\mu_2}{\int_{\mathcal{G}} f^2 d\mu_1}$$

<sup>&</sup>lt;sup>1</sup>We remark that a notion of a graph with boundary appears, at least implicitly, in [Dod84]. However, the treatment there deals with boundary edges of length one, in our terminology. For dealing with nodal regions it is essential to allow boundary edges of fractional length.

(see [CH53] or [Gar66] for a discussion of the classical situation). A minizing sequence for  $\mathcal{R}$  is any sequence of functions  $f_1, f_2, \cdots$  such that  $f_i$  acheives the infimum  $\mathcal{R}$  value among all nonzero  $f \in H$  which are  $\mu_1$  orthogonal to  $f_1, \ldots, f_{i-1}$ . The  $f_i$ 's are also called eigenfunctions, and the resulting values  $\lambda_i = \mathcal{R}(f_i)$  are called the eigenvalues of  $\mathcal{R}$ . We shall soon see that minimizing sequences exists; given this, it is easy to check that  $\lambda_i$  is independent of the choice of minizing sequence, as is the subspace  $E_{\lambda} \in H$  spanned by  $f_i$  such that  $\mathcal{R}(f_i) = \lambda$ . What differs here from the classical situation is that  $\mu_1$  is supported on n = |V| points, and so it only makes sense to speak of the first n eigenvalues and eigenfunctions. If f is non-zero and  $\mu_1$  orthogonal to the first n term of a minimizing sequence, then f vanishes on V and so, in some sense,  $\mathcal{R}(f) = +\infty$ ; one can take this as a definition, e.g. so that  $\lambda_i = +\infty$  for i > n, but this is of little concern to us.

By an edgewise linear function we mean a function  $f \in H$  whose restriction to each edge is a linear function.

**Proposition 2.1** Among all  $f \in H$  with given values at the vertices,  $\mathcal{R}(f)$  is minized at and only at the edgewise linear function with those vertex values. In particular the eigenvalues and eigenfunctions of  $\mathcal{R}$  exist (i.e. minimizing sequences exist) and are those of the Laplacian of G (i.e. the restriction of  $\mathcal{R}$  eigenfunctions to V are the Laplacian eigenvectors).

**Proof** The first statement follows from the easy fact that  $\int_0^1 (f')^2 dx$  subject to f(0) = a, f(1) = b is minized (over any reasonable class of functions) precisely when f is linear. For edgewise linear functions  $\mathcal{R}$  and  $\mathcal{R}_{\Delta}$  agree, and the second statement follows from linear algebra (using the self-adjointness of the Laplacian on  $L^2(V)$ ).

The nodal regions of an eigenfunction f of  $\mathcal{R}$  are the connected components of  $f^{-1}(\mathbf{R} - \{0\})$ . According to the classical theory f restricted to any of its nodal regions should be the first eigenvalue of the nodal region. So we must define what a graph with boundary should mean.

By a graph with boundary we mean an undirected graph  $G = (V \cup \partial V, E \cup \partial E)$  along with a "length,"  $c_e \in (0,1]$ , for each edge in  $e \in \partial E$  such that each E edge has both endpoints in V and each  $\partial E$  edge has exactly one endpoint in V and one in  $\partial V$ . We refer to V and  $\partial V$  vertices as interior and boundary vertices, respectively; similarly for the edges. By the geometric realization,  $\mathcal{G}$ , of G we mean the metric space as before except that for each  $e \in \partial E$  edge we glue in an interval of length  $c_e$ . We define  $\mu_1$  to be the measure which counts points in V, and  $\mu_2$  as Lebesgue measure on  $E \cup \partial E$ . We define  $\mathcal{R}$  and H as before; we define  $H_0$  to be the subspaces of functions in H which vanish at  $\partial V$ . Notice that if  $e = \{v, v'\}$  is a  $\partial E$  edge with  $v \in V$ , then if f is linear along e we have

$$\int_{e} |\nabla f|^{2} d\mu_{2} = \int_{e} (f(v) - f(v'))^{2} / c_{e}^{2} d\mu_{2} = (f(v) - f(v'))^{2} / c_{e}.$$

For this reason we set  $\Delta_0 = D_0 - A$  where A is the adjacency matrix restricted to V and where  $D_0$  is the diagonal matrix whose entry corresponding to  $v \in V$  is

$$(D_0)_{v,v} = \sum_{e \ni v} \frac{1}{c_e} \tag{2.1}$$

with  $c_e$  taken to be 1 for  $e \in E$ . An easy calculation shows that the associated Rayleigh quotient,  $\mathcal{R}_{\Delta_0}$ , agrees with  $\mathcal{R}$  on the set of edgewise linear functions in  $H_0$ . Again, minizing sequences for  $\mathcal{R}$  restricted to  $H_0$  exist and are edgewise linear; we call the associated eigenvalues and eigenfunctions the *Dirichlet* eigenvalues and eigenfunctions; they are those of  $\Delta_0$  (when restricted to V). In summary we have

The reader will note that the set  $\partial V$  plays no essential role in defining the Dirichlet eigenpairs; one can assume  $\partial V$  consists of one point, or that it has a distinct vertex for each  $\partial E$  edge, without affecting matters. In the latter case we say that G has separated boundary.

Given that G has separated boundary, we define the Neumann eigenvalues and eigenfunctions as those corresponding to  $\mathcal{R}$  on H. They exist, the eigenfunctions are constant along  $\partial E$  edges, and the eigenpairs are easily seen to be the same as those of  $\Delta_1 = D_1 - A$  where  $D_1$  is the diagonal matrix on V whose corresponding v entry is the number of neighbors it has in V (i.e.  $\Delta_1$  is just the usual Laplacian on V, ignoring all boundary edges and vertices). More generally we can define mixed boundary conditions; namely for each  $\partial E$  edge  $\{v_w, w\}$ , with  $w \in \partial V$  (we assume G is separated so that to each such w there corresponds only one v), we specify a boundary condition of one of the following forms: (1) Dirichlet, i.e. f(w) = 0, (2) Neumann, i.e. f(w) unrestricted, or (3)  $f(w) = a_w f(v_w)$  for some constant  $a_w \geq 0$ . Again it is easy to see that for this subspace H' of H eigenpairs exist and agree with a Laplacian  $\widetilde{\Delta} = \widetilde{D} - A$  for  $\widetilde{D}$  the diagonal matrix whose entry at v is

$$(\widetilde{D})_{v,v} = \sum_{\{w \mid (v,w) \in \partial E\}} \frac{|a_w - 1|}{c_{(v,w)}} + \sum_{\{e \in E \mid v \in e\}} 1$$

(with  $a_w$  being 0,1 for Dirichlet and Neumann conditions respectively). The Neumann condition is the same as requiring  $f(w) = f(v_w)$ , for this choice of f(w) minimizes the numerator of  $\mathcal{R}$  without affecting the denominator. In the above we require  $a_w \geq 0$  so as to have f(w) and  $f(v_w)$  of the same sign; if not, we would have to define the notion of nodal region more carefully. In any case there is no essential loss of generality in assuming  $a_w \geq 0$ , for  $\widetilde{D}$  depends only on the value of  $|a_w - 1|$ . In summary, we have shown:

**Proposition 2.2** The Dirichlet eigenpairs of a graph with boundary, G, i.e. successive orthogonal minimizers of the  $\mathcal{R}$  restricted to functions vanishing on the boundary of G, are the same as the eigenpairs of the matrix on interior vertices  $\Delta_0 = D_0 - A$  as above, whose eigenfunctions are extended by edgewise linearity to functions on all of G. Similarly for Neumann and mixed boundary conditions on graphs with separated boundary, with  $\Delta_1$  and  $\widetilde{\Delta}$  defined as above.

If  $G_1$  and  $G_2$  are graphs with separated boundary, we say that  $G_2$  is an extension of  $G_1$ , written  $G_1 \subseteq G_2$ , if there exists an isometric imbedding (i.e. 1 to 1) of the realization of  $G_1$  into  $G_2$  which preserves the degree of each interior vertex. In other words,  $G_1$  is obtained from  $G_2$  by declaring some of its interior vertices to be boundary vertices and by shortening some of the boundary edges' lengths (and ignoring any edges now connecting two boundary vertices). We also say that  $G_1$  is contained in  $G_2$ . If  $G_1$  and  $G_2$  are connected graphs and the above embedding is not onto we say that  $G_2$  is a strict extension of  $G_1$ , written  $G_1 \subset G_2$ . The

concept of being a strict extension implies both graphs involved are connected. In the above situations, any boundary conditions on  $G_2$  induces boundary conditions on  $G_1$  by imposing Dirichlet conditions on all exterior vertices of  $G_1$  which are interior to  $G_2$ , and keeping the same  $a_w$ 's on all exterior vertices, w, of  $G_1$  which are also exterior vertices of  $G_2$ ; we call these the *Dirichlet induced* boundary conditions. We similarly define the *Neumann induced* boundary conditions.

We say that  $G_1 \subset G_2$  freely if  $G_1 \subset G_2$  with  $G_2 - G_1$  having no cycles (i.e. is simply connected, i.e. every two points have at most one path joining them).

Given an eigenfunction of a graph, G, with any boundary conditions, each of its nodal regions determines a graph with boundary, G'; when there is no confusion we will identify a graph (possibly with boundary) with its realization.

At this point all of the classical theorems about monotonicity of eigenvalues as functions of the domain and of nodal regions of eigenfunctions go through essentially verbatim. In the below "eigenfunction" means an eigenfunction with any boundary conditions, unless otherwise specified.

**Theorem 2.3** The min-max and max-min principles hold for R acting on H or any subspace of H determined by some boulary conditions as mentioned above. The first eigenfunction of a connected graph has a strict sign on its interior vertices, i.e. all values on interior vertices are positive or all are negative; the first eigenvalue has multiplicity one, i.e. the first eigenfunction is uniquely determined up to scalar multiple; all higher eigenfunctions have both positive and negative values on interior vertices. Any nodal region  $G' \subset G$  of an eigenfunction, f, on G of eigenvalue  $\lambda$ , has  $\lambda$  as its first eigenvalue with boundary conditions on G' being those Dirichlet induced from G, and the restriction of f is the corresponding eigenfunction. If  $G_1 \subseteq G_2$  then the k-th eigenvalue of  $G_2$  for given boundary conditions is  $\leq$  that of  $G_1$  with the Dirichlet induced boundary conditions; in particular, this holds for the Dirichlet eigenvalues of  $G_1, G_2$ . Same with  $\subset$  and < for the first eigenvalue for any boundary conditions on  $G_2$  except the Neumann condition (on all exterior vertices). The k-th Dirichlet eigenvalue is > the k-th Neumann eigenvalue of a graph, with equality for the first eigenvalue iff the graph has no boundary. Any k-th eigenfunction of a connected graph has at most k nodal regions. If some of the quantities  $|a_w-1|$  in the mixed boundary conditions are increased, the eigenvalues cannot decrease; similarly if edges are added to the graph, and similarly if boundary edge lengths are decreased. If  $G_1, \ldots, G_r$  are contained in G with disjoint interiors (i.e. intersecting only on boundary vertices), then the k-th smallest eigenvalue among all the  $G_i$  Dirichlet eigenvalues is  $\geq$  the k-th Dirichlet eigenvalue of  $G_i$ ; if the  $G_i$  exhaust G, then the reverse inequality holds for Neumann eigenvalues; the same holds for any boundary conditions on G, with respectively Dirichlet and Neumann induced boundary conditions on the  $G_i$ .

**Proof** This follows verbatim from the classical theory (see, say, [CH53] or [Gar66]). In any case all statements are easy, and we include their proof in the appendix.

We remark that the situation with graphs can exhibit some phenomenon not present in analysis. For one thing, eigenfunctions in a connected graph can vanish in balls of arbitrarily large radius (a tree of degree  $\geq 3$  at all interior vertices has eigenfunctions which are zero at all interior vertices and assumes values -1,1,0 at the leaves); in analysis all the eigenfunctions are necessarily analytic (see, for example, [GT83, Fri69]) and hence can't vanish in any open set.

Similarly, if  $G_1 \subset G_2$  it is not true that all the Dirichlet eigenvalues need strictly increase (which is true in analysis, again by analyticity, assuming  $G_2 - G_1$  has nonempty interior); similarly many of the theorems of strict inequality of eigenvalues don't apply in general to graphs.

For an example of this phenomenon, consider an "hour glass" type graph, G: fix two isomorphic connected graphs,  $H_1, H_2$ , and fix  $v_i \in H_i$  which correspond to each other under the isomorphism. Let G consist of the  $H_i$  and an extra vertex v and the two extra edges  $(v, v_1)$  and  $(v, v_2)$ . Denote by  $H'_i$  the graph  $H_i$  union the boundary edge  $(v_i, v)$ ; and denote its first Dirichlet eigenpair  $f, \nu$ . We claim that the second Neumann eigenvalue of G is  $\nu$ . Clearly the function which is f on  $H'_1$  and -f on  $H'_2$  is orthogonal to the first eigenfunction and achieves a Rayleigh quotient of  $\nu$ . On the other hand, any eigenfunction with both positive and negative values has one of its nodal regions is contained strictly within one of the  $H'_i$ , having the same boundary conditions as those Dirichlet induced from  $H'_i$  (this follows from considering the eigenfunction's value at v); thus no other second eigenfunction can have eigenvalue smaller than  $\nu$ . On the other hand, if form G' by adding a vertex w and the edge (v, w), then the same argument shows that  $\nu$  is also the second Neumann eigenvalue of G' (with eigenfunction as before and extended by zero to w). Hence a graph can strictly increase without the same being true for all its (Neumann, Dirichlet, etc.) eigenvalues.

In classical graph theory there is an "interlacing theorem" for the eigenvalues of the adjacency matrix (or Laplacian) of a graph, G, and any graph, G', obtained by removing a vertex, v. That such a theorem holds for the Laplacian is contained in the above (half of this follows from the statement about Dirichlet eigenvalues when  $G' \subseteq G$ , the other half follows from the last statement about comparing Neumann eigenvalues, taking  $G_1$  to be G with all edges to v being replaced by separate boundary edges, and taking  $G_2$  to consist of v).

The theorem that each k-th eigenfunction has at most k nodal regions is known as the Hilbert nodal region theorem. Some of the above results, including this one, can be stated (and proven) in elementary classical terminology of graphs—e.g. the induced subgraph on the vertex subset of positive eigenvector values, and that on negative values, have at most k connected components in total. However, for us it is important that the eigenfunction/value restricted to a nodal region is a first eigenfunction/value of that region, and for this the notion of a graph with boundary is essential.

Also, as in the classical case we can vary the boundary of G slightly and analyze the change in the eigenvalues and eigenfunctions. Similarly if we change, say, the coefficients of  $\mu_2$  or  $\mu_1$  (meaning we are allowed to weight the vertices differently). In analysis this is needed to prove more subtle theorems, such as what happens to Neumann eigenvalues when you change the domain. In the case of graphs such techniques are not needed for any theorem known to the author. For example, an easy argument gives:

**Theorem 2.4** If  $G_1 \subseteq G_2$  freely, then the k-th Neumann eigenvalue of  $G_1$  is at least that of  $G_2$ .

**Proof** If  $G_1 \subseteq G_2$  freely, and  $G_2$  has exactly l more interior vertices than  $G_1$ , then their Laplacian matrices (i.e.  $G_1$ 's extended by zero to the additional interior vertices in  $G_2$ ) differ by a matrix, M, which is the sum of l matrices, each vanishing on all but two columns and rows where it is of the form

 $\left[\begin{array}{cc} 1 & -1 \\ -1 & 1 \end{array}\right].$ 

Hence M is a positive semi-definite matrix of rank  $\leq l$  (actually = l in this case), and the claim follows from the max-min principle.

Note that if  $G_1 \subseteq G_2$  but not necessarily freely, then some of the Neumann eigenvalues of  $G_2$  can be smaller than those of  $G_1$ ; indeed, if  $G_1$  is any d-regular graph (without boundary) and  $G_2$  is obtained by adding a vertex, v, with an edge to every  $G_1$  vertex, then each higher Neumann eigenfunction of  $G_1$ ,  $f_2$ ,  $f_3$ ,  $\cdots$  extends by zero on v to an eigenfunction on  $G_2$  with eigenvalue = 1+ the old eigenvalue. Taking just about any d-regular graph, such as a large cycle, then gives an example of non-monotonicity in this situation.

One does not need to insist on separated boundaries to define Neumann and other boundary conditions. Indeed one can define Laplacians, and much of the above theory goes through. On the other hand, as the example in the previous paragraph suggests, monotonicity theorems for such boundary conditions often require free containment, and non-free containment can enlarge the eigenvalues (much as does adding edges to a graph).

We also remark that in analysis, the Neumann eigenvalues can increase when increasing the size of the domain, even in  $\mathbb{R}^2$ , as shown by an example such as that of [CH53], chapter VI, section 2.6. Namely, if two fixed squares are joined by a very thin rectangle, then as the thinness tends to zero the second Neumann eigenvalue tends to zero; however, this family of domains are contained in some large fixed rectangle, whose second Neumann eigenvalue is > 0.

The technique of comparing two graphs by a smoothly varying family of graphs from one to the other seems quite interesting and yields another proof of the above theorem. Furthermore, we will need the version of it which holds for the Dirichlet boundary conditions later in this paper.

**Theorem 2.5** Let G be a graph with boundary, having Dirichlet eigenvalues  $\nu_1 \leq \cdots \leq \nu_n$  with corresponding eigenfunctions  $f_1, \ldots, f_n$ . Let  $\sigma$  be any real valued function on the boundary edges, and for  $\epsilon > 0$  let  $G_{\epsilon}$  be the graph obtained from G by adding  $\epsilon \sigma(e)$  to the length of e for each boundary edge, e (we therefore assume  $\sigma$  is non-negative if  $c_e = 0$ , non-positive if  $c_e = 1$ ). Then for all  $i \leq n$  we have  $\nu_i(G_{\epsilon})$  is differentiable in  $\epsilon$  with

$$\left. \frac{\partial \nu_i(G_\epsilon)}{\partial \epsilon} \right|_{\epsilon=0} = \sum_{e \in \partial E} \sigma(e) \left( \frac{\partial f}{\partial \eta}(e) \right)^2,$$

where  $\partial/\partial\eta$  denotes the normal derivative of f at e, i.e. in the direction from interior to exterior vertex.

The above formula, with the summation over  $\partial E$  replaced by integration over the boudnary of a domain and with  $\partial/\partial\eta$  interpreted as the normal derivative is just the classical variational formula for Dirichlet eigenvalues.

**Proof** Perturbation theory tells us that if a symmetric matrix,  $M(\epsilon) = M_0 + \epsilon M_1 + O(\epsilon^2)$  has entries which are analytic functions in  $\epsilon$  near 0, then the *i*-th eigenvalue  $\nu_i$  of M is analytic in  $\epsilon$  near 0 and

 $\frac{d}{d\epsilon}\nu_i = (M_1 v_i, v_i)$ 

with  $v_i$  the *i*-th eigenfunction normalized to be of unit length. The  $\nu_i(G_{\epsilon})$ 's are given as the eigenvalues of the matrix  $\Delta_0(\epsilon) = D_0(\epsilon) - A$ , with  $D_0(\epsilon)$  given in equation 2.1, with all  $c_e$ 's being linear functions of  $\epsilon$ . It follows that the  $\nu_i$ 's are locally analytic in  $\epsilon$ . It suffices to verify the above formula for  $\sigma$  taking only one non-zero value, and this is a simple calcultion based on the above equation.

The above nodal region theory also allows us to generalize certain theorems from graph theory. For example, for a d-regular graph without boundary, G,  $\lambda_2(G)$  is related to its second Neumann eigenvalue,  $\nu_2$ , via  $\lambda_2 = d - \nu_2$ . Dodziuk, in [Dod84], proves a Cheeger type inequality for the first Dirichlet eigenvalue of a d-regular graph with, in our terminlogy, all boundary edges of length one. While the proof given there does not directly apply to  $\nu_2$  or  $\lambda_2$ , the theorem still applies to  $\lambda_2$  via the above theory. Namely, if G' is the nodal region of the second Neumann eigenfunction with the least number of vertices, then the graph  $\widetilde{G}$  obtained grom G' by making all boundary edges of length one has first Dirichlet eigenvalue no greater than  $\nu_2$ ; on the other hand Dodziuk's result applies to  $\widetilde{G}$ , and thus yields a Cheeger type inequality for  $\lambda_2$  in terms of the "magnification" constant (as in [Alo86]). This gives a proof of a Cheeger type inequality for  $\lambda_2$  which is very similar to the proof in analysis; Alon's theorem, in [Alo86], yields a better bound on  $\lambda_2$ , but his proof is trickier.

## 3 Eigenvalue Upper/Lower Bounds and Capacity

We restate part of theorem 2.3, which will be used to obtain lower bounds on  $\lambda_2$ . In this section we apply this only to graphs with boundary of length 1, where it is also follows from elementary linear algebra. This trick was exploited by Weyl in his proof of the growth rate of eigenvalues of the Laplacian.

**Proposition 3.1** Let  $G_1, \ldots, G_s$  be graphs with boundary which are disjointly (i.e. with disjoint interiors) contained in G. Then the number of Dirichlet eigenvalues of G which are  $\leq \lambda$  for any  $\lambda$  is at least the total number of Dirichlet eigenvalues of the  $G_i$  which are  $\leq \lambda$ .

Let  $T_{d,k}$  be the d-regular undirected tree of depth k; so  $T_{d,k}$  has a root and  $d(d-1)^{i-1}$  vertices at distance i for each  $i \leq k$ .

Proposition 3.2 The largest eigenvalue of  $T_{d,k}$ 's adjacency matrix is

$$\lambda_1^{d,k} = 2\sqrt{d-1}\cos\theta_{d,k},$$

where  $\theta_{d,k}$  is the smallest positive solution,  $\theta$ , of  $g_k(\theta) = d/(2d-2)$ , where

$$g_k(\theta) \equiv \frac{\sin((k+1)\theta)\cos(\theta)}{\sin(k\theta)}.$$

We have  $\theta_{2,k} = \pi/(2k+2)$ , and for all  $d \geq 3$  we have  $\theta_{d,k} \in [\pi/(k+5), \pi/(k+1)]$ .

**Proof** For  $x \in T_k$  let  $\rho(x)$  denote x's distance to the root. The function

$$f(x) = (d-1)^{-\rho(x)/2} \sin\left(\pi - (k+1-\rho(x))\theta\right)$$
 (3.1)

is positive on  $T_k$  provided that  $\theta>0$  and  $\pi-(k+1)\theta>0$ . Away from the root we have  $Af=\lambda f$  with A the adjacency matrix and  $\lambda=2\sqrt{d-1}\cos\theta_{d,k}$ ; at the root we have  $Af=\lambda f$  provided that  $g_k(\theta)=d/(2d-2)$ . Since  $g_k(\pi/(2k+2))=1$  and  $g_k(\pi/(k+1))=0$  (and  $g_k'(\theta)<0$  on  $[\pi/(2k+2),\pi/(k+1)]$ ), a  $\theta$  making f a positive eigenfunction exists, and for this  $\theta$  the Perron-Frobenius theorem (see the appendix) implies that the largest eigenvalue of  $T_k$  is f's eigenvalue. An analysis of  $h(k)=g_k(\pi/(k+5))$  shows that  $h\geq 4/5$  for all  $k\geq 1$ ; indeed, setting s(x)=h(1/x) we check that s(0)=4/5 and s'(x) has the same sign as  $4\sin(\alpha)-\sin(4\alpha)$  with  $\alpha=2\pi x/(1+5x)$ ; since  $\alpha\leq\pi/2$  for all  $x\geq 0$ , s is monotone increasing there, and so s(x)>4/5 for x>0.

For a non-negative symmetric square matrix (directed graph), its (average) valence<sup>2</sup> is its largest positive eigenvalue (that of its adjacency matrix). The information theory literature usually speaks of the capacity, which is the log<sub>2</sub> of the valence.

We say that a graph contains a d-regular ball of radius k if there exists a vertex  $v_0$  such that every vertex of distance < k has degree d. We note that among all such graphs,  $T_{d,k}$  has the smallest capacity.

**Lemma 3.3** Let G contain a d-regular ball of radius k about a vertex  $v_0$ . Then the capacity of G is at least that of  $T_{d,k}$  (with equality iff G is isomorphic to  $T_{d,k}$ ).

**Proof** Clearly there exists a local isomorphism of  $T_{d,k}$  into G; fix one such, say  $\phi$ . Via  $\phi$ , every walk from the root, r, of  $T_{d,k}$  to itself gives rise to a walk in G from  $v_0$  to itself of the same length, and this map is one to one. The capacity inequality follows by Shannon's algorithm (see the appendix) or by considering traces of high powers of the adjacency matrices; strict inequality when G is not  $T_{d,k}$  follows from Shannon's algorithm.

<sup>&</sup>lt;sup>2</sup>We adopt the terminology in [AFKM86], except that we will almost always omit the word average. In the information theory literature the word *capacity* is reserved for the log<sub>2</sub> of this quantity.

**Definition 3.4** A graph with boundary is d-regular if every interior vertex has degree d. We then define its (average) valence to be  $d - \nu_1$ , where  $\nu_1$  is its first Dirichlet eigenvalue.

Since the first Dirichlet eigenfunction is positive on the vertices,  $d - \nu_1$  is also the average valence (see the appendix) of the matrix

$$dI - \Delta_0(G) = A_0(G) = A(G) - L(G)$$

where I is the identity matrix,  $\Delta_0$  is as in section 2, A is the adjacency matrix restricted to V, and L is the diagonal matrix whose diagonal entry at v is

$$\sum_{e\ni v} \left(\frac{1}{c_e} - 1\right).$$

(In the above the sum can be taken over only  $\partial E$  edges since the summand vanishes at E edges.) In particular we see that the usual notion of valence of an adjacency matrix coincides with this definition in that:

**Proposition 3.5** If all of the  $\partial E$  edges have length 1, then L=0 and the valence of G is just the valence of A.

Corollary 3.6 Let G be a d-regular graph with a subset of r points each of distance  $\geq 2k$  from one another. Then  $\lambda_r(G) \geq \lambda_1^{d,k}$ .

**Corollary 3.7** For a d-regular graph on n vertices with  $d \geq 3$ , the r+1-th largest eigenvalue is at least

$$2\sqrt{d-1}\left(1-\frac{\pi^2}{2k^2}+O\left(k^{-4}\right)\right).$$

where  $k = (\log_{d-1}(n/r))/2$ .

We can give a weaker but elementary version of the above corollaries, by noting that the function which is zero for  $\rho(x) > k$  and otherwise equal to f(x) as in equation 3.1 with  $\theta = \pi/(2k+2) = \theta_{2,k}$  satisfies  $Af \geq \lambda f$  with  $\lambda = 2\sqrt{d-1}\cos\theta$ . If G has no odd length cycles of length  $\leq 2k$  then using the function  $(-1)^{\rho(x)}f(x)$  shows that:

**Proposition 3.8** If G is a d-regular graph with a subset of r points each of distance  $\geq 2k$  from one another, and contains no odd cycle of length  $\leq 2k$ , then  $\lambda_{n-r}(G) \leq -\lambda_1^{2,k} =$ 

$$2\sqrt{d-1}\left(1-\frac{2\pi^2}{k^2}+O(k^{-4})\right).$$

This also gives an easy proof of the last corollary with a slightly worse  $k^{-2}$  constant.

We remark that under very weak assumptions about the graph, much more is true. Namely, given a sequence,  $G_1, G_2, \ldots$  of d-regular graphs whose number of vertices,  $n_i = |V_i|$  goes to infinity, such that the number of cycles of length k is o(n) for every fixed k, then the eigenvalue distribution converges weakly to a distribution depending only on d (see [McK81], also [LPS88]).

For an arbitrary d-regular graph on n vertices,  $G_{n,d}$ , it is easy to see that  $\lambda_2 \geq \sqrt{d(n-d)/(n-1)}$  by considering a vertex and the n-d vertices to which it is not connected. The fact that  $\lim\inf f$  of  $\lambda_2(G_{n,d})$  tends to  $2\sqrt{d-1}$  for fixed d as  $n\to\infty$  appears in many places (implicitly in [McK81], perhaps earlier). The first explicit mention of of a lower bound on  $\lambda_2$  in term of d and n occurs in Alon's paper [Alo86], as to due himself and Boppana; later he gives a correct statement and a proof, in [Nil91], of a  $2\sqrt{d-1}(1-O(1/k))$  lower bound on  $\lambda_2$ . The fact that one can replace O(1/k) with  $O(1/k^2)$  seems to have been unnoticed. A. Lubotzky has pointed out to the author that the fact that the lim inf of  $\lambda_r$  tends to  $2\sqrt{d-1}$  for fixed r and d as  $n\to\infty$  has also been observed by M. Burger.

## 4 Towards a Faber-Krahn Type Inequality

To understand d-regular graphs better we wish to study isoperimetric type problems on the d-regular infinite tree,  $T_d$ . For one thing, every d-regular graph is a quotient of  $T_d$ . For another, any graph with no short cycles looks locally like a tree, and so results on  $T_d$  will translate into results for isoperimetric problems restricted to small sets in the graph. Futhermore, many graphs with good isoperimetric properties, such as being expanders, contain no short cycles (see, e.g., [Chu88, LPS88] and proposition 4.2); intuitively speaking, a graph which expands well should have no simple unforced relationships between its edges, and hence should contain no short cycles.

Our experience from Euclidean space (or discretized versions of it), in which the best sets with respect to almost any type of isoperimetric problems are balls, is certainly not true in  $T_d$ . Indeed, for a subset of vertices, A, of  $T_d$ , let  $\Gamma(A)$  denote set of vertices connected to A by an edge, and let E(A, A') be the set of edges with one endpoint in A and one in, A', the complement of A. It is typical to ask for an A of a given size, how small can  $|\Gamma(A) - A|$  or |E(A, A')| be and what are the sets A which achieve these values. It is easy to see that any connected set is an isoperimetrical extremal set in this sense.

Proposition 4.1 For any nonempty subset of vertices, A, of  $T_d$ , we have

$$|\Gamma(A) - A| = (d-2)(|A|-2) + 2(d-1) - \tilde{c}, \qquad |E(A,A')| = (d-2)(|A|-2) + 2(d-1) - 2c$$

with  $c, \tilde{c} \geq 0$ ; more precisely, c is the number of connected components of A (as an induced subgraph of  $T_d$ ) minus one, and  $\tilde{c} \geq c$  is 2c minus the number of pairs of connected components of A whose distance is 2.

As a corollary we get results such as:

**Proposition 4.2** In a d-regular graph whose shortest cycle is of length  $\ell \geq$ , respectively, 5 and 4, among all sets of size  $\ell$ , A, the ones minizing, respectively,  $|\Gamma(A) - A|$  and |E(A, A')|, are precisely those cycles of length  $\ell$ .

So in finite graphs without small cycles, good isoperimetric sets can be very "thin sets."

Closer to our intution about Euclidean space, we conjecture that a Faber-Krahn type inequality holds in  $T_d$ . Namely, for an open subset G of  $T_d$ , let  $\nu_1$ , the first Dirichlet eigenvalue of G, be the minimizer of the Rayleigh quotient among nonzero  $T_d$  functions vanishing outside of the interior of G (so that  $\nu_1 = \infty$  if G contains no vertices); this is the same as the first Dirichlet eigenvalue of G viewed as a graph with boundary, where all components of G not containing a vertex are discarded. In this way it is easy to see that  $\nu_1$  is continuous as a function of G in the metric  $\rho(G, G') = \mu_2(G - G') + \mu_2(G' - G)$  where  $\mu_2$  is the measure of section 2 (and where the range, the non-negative reals union  $\infty$ , is topologized as usual).

For any  $p \in T_d$ , not necessarily a vertex, and r > 0, the set of all points of distance  $\leq r$  from p,  $B_r(p)$ , determines a subgraph of  $T_d$  with boundary assuming  $B_r(p)$  contains at least one vertex.

Conjecture 4.3 Among all  $G \subset T_d$  with  $\mu_2(G) = S$  fixed, the first Dirichlet eigenvalue,  $\nu_1(G)$ , is minimized when and only when G is a ball centered at a vertex.

We remark that a graph with boundary, G, is isomorphic to a subgraph of  $T_d$  iff it is d-regular and is a tree.

At present we can only prove a weaker statement:

**Theorem 4.4** For any fixed S > 0, there is a  $G_S \subset T_d$  of  $\mu_2$ -measure S whose  $\nu_1$  attains the infimum over all  $\nu_1$ 's of graphs with  $\mu_2$ -measure S. Any such  $G_S$  is connected. Assuming  $d \geq 5$ , for any R there is an  $S_0$  such that any  $G_S$  with  $S \geq S_0$  contains a ball of radius  $\geq R$ .

**Proof** Let G be a subgraph of  $T_d$ . If G has more than one connected component, then  $\nu_1(G)$  is clearly the smallest  $\nu_1(G')$  ranging over all connected components, G', of G. By theorem 2.3,  $\nu_1(G)$  would strictly decrease if we replaced G by any connected subgraph of  $T_d$  strinctly containing G'; such as subgraph can be chosen to have measure  $g = \mu_2(G)$ .

So consider a sequence  $G_1, G_2, \ldots$  of graphs of  $\mu_2$ -measure S whose  $\nu_1$  tends to the infimum. We can assume the  $G_i$  are connected, by the above.

Consider all pairs (G, p) of a connected subgraph of  $T_d$ , G, of  $\mu_2$ -measure S with a point  $p \in G$ . We can specify all pairs (G, p), up to isomorphism, via a finite set of bounded real coordinates, with  $\nu_1$  being a continuous function of these coordinates (fix for each G an embedding into  $T_d$ , and for each "non-backtracking" path of length S from p in  $T_d$  define a coordinate on G by its length along that path). It follows from compactness that there exists a limit point, G, of  $G_1, G_2, \ldots$ , with  $\mu_2(G) = S$ , and by continuity G attains the infimum of  $\nu_1$ .

Now fix any G of  $\mu_2$ -measure S minizing  $\nu_1$  (which is necessarily connected by the above), and let f be a corresponding eigenfunction. Theorem 2.5 implies that the normal derivative of f at all boundary edges is the same. Hence we can take this to be 1 by rescaling f (this means that f is positive in the interior of G). This fact alone gives us some symmetries in

G; namely, all boundary edges to a fixed vertex v necessarily have the same length,  $c_v$ , and furthermore  $f(v) = c_v$ .

To obtain more symmetries it is helpful to note:

**Lemma 4.5** There is no geodesic path of vertices  $p = (v_1, v_2, ..., v_r)$  in G such that for some i < j < k we have  $f(v_j)$  is less than both  $f(v_i)$  and  $f(v_k)$ .

Corollary 4.6 Let w be a maximum of f on G. Then along any geodesic from w to a boundary vertex, f is non-increasing.

Here we say a path is *geodesic* if it is the shortest path joining its endpoints (this path is always unique in a tree).

**Proof** For the lemma, first of all, we may assume by extending p that  $v_1$  and  $v_r$  are boundary vertices. Notice that since G is a subgraph of  $T_d$ , G consists of the path p plus trees  $T_i$  rooted at  $v_i$  for  $2 \le i \le r-1$ , each of degree d-2. Given a permutation,  $\sigma$ , of  $\{2,\ldots,r-1\}$  we can form a new graph  $G_{\sigma}$  consisting of the path  $p_{\sigma} = (v_1, v_{\sigma(2)}, \ldots, v_{\sigma(r-1)}, v_r)$  and attaching  $T_i$  to  $v_i$  as before;  $G_{\sigma}$  is a d-regular graph with boundary of the same  $\mu_2$  measure as G. We can define  $f_{\sigma}$  as the edgewise linear function coming from the values of f at the vertices in the obvious way;  $\nu_1$  of  $G_{\sigma}$  will certainly be less than that of G if the integral of  $|\nabla f_{\sigma}|^2$  along  $p_{\sigma}$  is less than than of  $|\nabla f|^2$  along p (since the  $\mu_1$  integral of  $f^2$  and  $f^2_{\sigma}$  are clearly equal). It suffices to show that such a  $\sigma$  exists under the assumption of the lemma.

So let  $m_1, m_2$  denote the respective maxima of f from  $v_1$  to  $v_j$  and from  $v_j$  to  $v_r$ . We can assume  $m_1 \leq m_2$ . Let  $i_1 < i_2 < j$  be two integers such that  $f(v_n) = m_1$  for  $i_1 \leq n \leq i_2$  and  $f(v_n) < m_1$  for  $n = i_1 - 1$  and  $n = i_2 + 1$ . Let  $k_1 > j$  be any integer such that  $f(v_{k_1}) \leq m_1 \leq f(v_{k_1+1})$ . Let  $\sigma$  be the permutation formed by removing the interval  $[i_1, i_2]$  and pasting it in between  $k_1$  and  $k_1 + 1$ . Then an easy calculation shows that  $\sigma$  is of the desired type, i.e.  $\nu_1(G_\sigma) < \nu_1(G)$ ; indeed, if  $\epsilon_1 = m_1 - f(v_{i_1-1})$ ,  $\epsilon_2 = m_1 - f(v_{i_2+1})$ ,  $\epsilon_3 = m_1 - f(v_{k_1})$ ,  $\epsilon_4 = f(v_{k_1+1}) - m_1$ , then  $\epsilon_i \geq 0$  and  $\epsilon_1, \epsilon_2 > 0$ , and the calculation boils down to the fact that

$$\epsilon_1^2 + \epsilon_2^2 + (\epsilon_3 + \epsilon_4)^2 > (\epsilon_2 - \epsilon_1)^2 + \epsilon_3^2 + \epsilon_4^2.$$

The corollary follows immediately from the lemma.

Our general strategy is as follows. Let the maximum value of f occur at the vertex v. View v as the root of  $T_d$ , and consider the heights of all the other vertices of G (i.e. their distance to v). Let  $V_1$  be the set of "maximal height" interior vertices, i.e. interior vertices such that all of their edges away from v are boundary edges. We know that for each  $v \in V_1$ , the length of its d-1 boundary edges are the same, and if this value is  $c_v$  then also  $f(v) = c_v$ . This then determines the value of f at the parent of v, in terms of v0. We will be able to show that every child of this parent, like v1, is an interior node with v2 boundary edges each of the same length as those of v3. Our strategy is to prove that as we move up the

tree from  $V_1$ , we will be able to claim that all interior vertices have balanced d-1 regular trees ascending from them. We will be able to do this up to some given length, depending on the value of  $\nu_1$ .

We make this description more precise as follows. For interior vertices v, w of G joined by an edge e, we say that v spawns a balanced tree at w of type (k, c) if all geodesics from w to boundary vertices through v are of length k+c, where k is an integer and  $c \in (0,1]$ . This means that the connected component, T, of v in G-e, is a balanced (d-1)-ary tree rooted at v in the usual sense, with all boundary edges of length c. The above shows that if  $w \in V_2$  is connected to a  $v \in V_1$  with all other neighbors of v being boundary vertices, then the tree spawned by v from w (in the usual sense) is balanced in the above sense. In this case the value of f, being an eigenfunction of the matrix  $D_0 - A$  (as in equation 2.1), must satisfy

$$(1 + (d-1)/c)f(v_1) - f(v_2) = \nu f(v_1),$$

i.e.

$$f(v_2) = c + (d-1) - \nu c.$$

Similarly if v spawns a tree of type (k-1,c) from w, then f(w) is given as  $f_k = f_k(c)$  determined by the recurrence relation

$$f_i = (d - \nu)f_{i-1} - (d-1)f_{i-2} \quad \forall i \ge 3, \qquad f_1 = c, \ f_2 = (d-1) + (1-\nu)c.$$

This can be used to derive some weak symmetry properties of G. Notice that by theorem 3.2 and theorem 2.3 we know that  $\nu > d - 2\sqrt{d-1}$ .

**Lemma 4.7** If  $v_1, v_2$  spawn balanced trees of type  $(k_1, c_1)$  and  $(k_2, c_2)$  from the same vertex w, then  $f_{k_1}(c_1) = f_{k_2}(c_2)$ . In particular, if  $d \ge 5$  and  $\nu = d - (2 + \epsilon)\sqrt{d-1}$ , then either (1)  $k_1 = k_2$  and  $c_1 = c_2$ , or (2) one of  $k_1, k_2$  is  $\ge R(\epsilon)$ , with R a function of  $\epsilon$  (independent of d) which  $\to \infty$  as  $\epsilon \to 0$ .

**Proof** The first statement is obvious. For the second part, note that (by induction on k) the functions  $f_k(c)$  are linear functions of c. Denote  $f_k(0)$ ,  $f_k(1)$  by  $\alpha_k, \beta_k$ . We have  $f_k(c) = (1-c)\alpha_k + c\beta_k$ . Let  $I_k$  be the interval with endpoints  $\alpha_k, \beta_k$ . It suffices to show that  $I_1, \dots, I_R$  are disjoint intervals none of which is a point (i.e.  $\alpha_k \neq \beta_k$  for  $1 \leq k \leq R$ ), for  $R = R(\epsilon)$  as in the statement of the lemma.

Notice that  $\alpha_1 = 0$ ,  $\alpha_2 = d - 1$ ,  $\alpha_3 = (d - 1)(d - \nu)$ , and  $\beta_1 = 1$ ,  $\beta_2 = d - \nu$ . It then follows that  $\alpha_{k+1} = (d-1)\beta_k$  for  $k \ge 1$ .

Let  $r_k = \beta_k/\beta_{k-1}$  for  $k \geq 2$ . Now  $r_2 = d - \nu < d - 1$  since  $d \geq 5$ , and  $r_k$  satisfies

$$r_k = (d - \nu) - \frac{d-1}{r_{k-1}} \quad \forall k \ge 3.$$

It easily follows that  $r_k$  is decreasing in k while  $r_k$  is positive, and that the first k time  $r_k$  becomes nonpositive,  $r_{k+1}$  is positive (or undefined, if  $r_k = 0$ ). In other words,  $r_k$  is decreasing at least until the first time when both  $\beta_k$  and  $\alpha_k$  are non-positive. Since the first

Dirichlet eigenfunction is positive at interior nodes,  $f_k(c)$  must be postive for all  $k \leq K$  where K is the larger of  $k_1, k_2$ , which implies that  $d-1 > r_2 > \cdots > r_K$ .

In particular we have  $\alpha_k > \beta_k$  for  $k \geq 2$  and  $I_1 = [0,1]$ , so none of the intervals  $I_k$ ,  $k \leq K$ , is a point. It suffices to show that  $\beta_k > \alpha_{k-1}$  for  $k \leq R(\epsilon)$  to show that  $I_1, \ldots, I_R$  are disjoint. Setting  $s_k = r_k/\sqrt{d-1}$ , this reduces to showing that the sequence given by  $s_2 = 2 + \epsilon$ ,

$$s_k = (2 + \epsilon) - \frac{1}{s_{k-1}} \quad \forall k \ge 3$$

has  $s_3 > \cdots > s_R > 1$  for  $R = R(\epsilon)$ . However at  $\epsilon = 0$  the  $s_k$  are just the Newton iterates for the equation  $(x-1)^2 = 0$  at initial point x = 1, and so  $s_R > 1$ ; hence for sufficiently small  $\epsilon$  we have  $s_R(\epsilon) > 1$  as well.

The theorem follows at once, noting that for  $\nu_1(G_S)$  tends to  $d-2\sqrt{d-1}$  as  $S\to\infty$  by theorem 3.2 and theorem 2.3.

We finish the section by noting that the same claim about  $f_{k_1}(c_1) = f_{k_2}(c_2)$  can be derived from Shannon's algorithm. We outline how that calculation is performed.

Consider for  $B \in \mathbf{R}$  the matrix  $M_B \equiv BI - \Delta_0 = BI + A - D_0$  where I is the identity matrix; for sufficiently large B this matrix will have all positive entries. Its largest eigenvalue is  $B - \nu$ . Assume v spawns a balanced tree from w of type (k, c). By Shannon's algorithm,  $z_0 = 1/(B - \nu)$  is the smallest positive root of the equation  $P_w^G(z) = 1$ . We may write  $P_w^G$  as the sum of  $P_{w,v} + P_0$ , with  $P_{w,v}$  representing the contribution from the walks through v, and  $v_0$  representing the other walks. By convention, if  $v_0$  is a boundary vertex connected to  $v_0$  by an edge of length  $v_0$ , we say  $v_0$  spawns a balanced tree of type  $v_0$  from  $v_0$ .

**Lemma 4.8**  $P_{w,v}$  is given by  $\phi_{k_i}(c_i, z)$ , where the  $\phi_k$ 's are given by  $\phi_{-1}(c) = -z/c$ , and for  $i \geq 0$ ,

$$\phi_i(c,z) = z^2 \frac{1}{1 - m_i(c,z)},$$

where the  $m_i$  are given by

$$m_0(c,z) = \left( (B-1) - \frac{(d-1)}{c} \right) z,$$

and for i > 1

$$m_i(c, z) = (B - d)z + (d - 1)\phi_{i-1}(c).$$

**Proof** Follows easily by induction on i.

Let  $v_1, v_2$  spawn trees from w as in lemma 4.7. As small change in  $c_1, c_2$  gives rise to a new graph, whose  $\mu_2$  measure is the same as the original graph as long as

$$(d-1)^{k_1}c_1 + (d-1)^{k_2}c_2$$

is preserved. The new graph cannot have a larger  $P_w^G(z_0)$ , for if so it would have a larger  $B-\nu$  and so a smaller  $\nu$ . Hence we have

$$(d-1)^{-k_1} \frac{\partial}{\partial c} \phi(c_1, z_0) = (d-1)^{-k_2} \frac{\partial}{\partial c} \phi(c_2, z_0).$$

We wish to calculate under what circumstances this can occur.

To simplify the calculation, it will suffice to take a first order approximation to the above for large B. More precisely, set  $t = 1/(B - \nu)$ , which means that  $z_0 = t$ , and consider the above equation for small t > 0. We have for  $i \ge 1$ ,

$$\frac{\partial}{\partial c}\phi_i = z^2 \frac{1}{(1 - m_i)^2} \frac{\partial}{\partial c} m_i,$$

and for  $i \geq 2$ ,

$$\frac{\partial}{\partial c}m_i = (d-1)\frac{\partial}{\partial c}\phi_{i-1}.$$

It then follows that:

**Lemma 4.9**  $m_i(c,t) = 1 - tK_i(c) + O(t^2)$  where  $K_i(c)$  is the function given by

$$K_i(c) = d - \nu - \frac{1}{K_{i-1}(c)} \quad \forall i \ge 1, \qquad K_0(c) = \frac{d-1}{c} + 1 - \nu.$$

For all  $j \leq i$ ,  $K_j(c)$  is necessarily non-negative. Furthermore, for all  $i \geq 0$ ,

$$\frac{\partial}{\partial c}\phi_i(c,t) = \frac{(d-1)^{1+i}}{c^2} \frac{1}{\left(K_0(c)\cdots K_i(c)\right)^2} t + O(t^2).$$

The formula remains valid for i = -1 (if we omit all the K's).

Now it is easy to see that the product  $cK_0(c)K_1(c)\cdots K_i(c)$  is just  $f_i(c)$ , yielding an alternate derivation of the fact that  $f_{k_1}(c_1) = f_{k_2}(c_2)$ .

## 5 Coverings and Eigenvalues

We hope to obtain spectral information about a graph by studying the spectrum of its covers, at least to some extent (see [Bro86] for examples in analysis). If H is a covering (see appendix) of a connected graph G, then  $\lambda_1(H) = \lambda_1(G)$ . While  $\lambda_2(H) \ge \lambda_2(G)$  can hold with equality for some H, it seems unlikely that it would hold with equality for a fixed G and many H's. We wish to discuss the following question:

Question 5.1 Let G be an undirected finite graph. Under what conditions can there exist connected covers of G,  $H_i$ , with size tending to infinity and with  $\lambda_2(H_i) = \lambda_2(G)$  for all i? Can one give geometric conditions on G and/or the sequence  $\{H_i\}$  to preclude this behavior?

At some point the author entertained the possibility that in any such situation one would have  $\lambda_2(H_i) > \lambda_2(G)$  for sufficiently large i, for the following reason. The theorem in the last section implies that for any d-regular graph with no small cycles, its subgraphs with smallest first Dirichlet eigenvalue among all subgraphs of their size contain large balls (depending on the size of the graph and the length of its shortest cycle). This leads us to believe that for large graphs of bounded degree (or at least for those with no short cycle), the nodal regions of the second eigenfunction should be of large inradius, i.e. should contain large balls.

On the other hand, the second eigenvector of G lifts to an eigenvector on  $H_i$  whose nodal regions' inradii are bounded by those of the nodal regions on G. It seems unlikely, at least for graphs,  $H_i$ , with good isoperimetric properties, that their nodal regions would look "thin" in the above sense.

Unfortunately one can give examples to show that sometimes one can have  $\lambda_2(H_i) = \lambda_2(G)$  for all i. Namely, let  $T_i$  be any collection of d-regular graphs for some d whose  $\lambda_2$ 's are below a fixed constant  $\lambda$ , and let S be a d'-regular with  $\lambda_2(S)d > d\lambda$  ( $T_i$  can be taken to be graphs such as the  $Y^{p,q}$ 's described below with p fixed, and S then chosen to be a sufficiently large cycle graph). Then setting  $G = S \otimes B_d$  and  $H_i = S \otimes T_i$  gives such an example, since spectrum of  $G_1 \otimes G_2$  is just the pointwise products of the spectra of  $G_1, G_2$ .

In the preceding example the graphs in question have quite bad expansion properties. It is conceivable that collections  $\{H_i\}$  with more typical geometric properties (i.e. like those of random graphs) never have a persistent  $\lambda_2$  as above.

P. Sarnak has pointed out to the author the connections between the above questions and number theory. Namely, in [Iha66] Ihara uses the Brandt-Eichler Zahlentheorie der Quaternionenalgebren to give various (p+1)-regular graphs with the property that its second eigenvalue is  $\leq 2\sqrt{p}$ . Furthermore, all second and lower eigenvalues of these graphs are related to those of Frobenius in characteristic p acting on (the Jacobian of) the modular curve  $Y_0(\ell)$  (see, e.g., [Sar]). In particular, the non-persistence of the second eigenvalue for these graphs would imply that no eigenvalue of Frobenius as above can be purely real and positive.

We recall a special case of these graphs which have a simple description and which we will use later for numerical experiments. Namely, recall the graphs,  $X^{p,m}$ , with p prime  $\equiv 1 \pmod 4$  and m > 1 and relatively prime to p from [LPS88] (also in [Mar87]; see [Bie89] for a complete discussion of general m), which are quotients of trees generated by quaternions of norm p. We also recall the graphs,  $Y^{p,q}$  for q prime  $\equiv 1 \pmod 4$ , which are quotients of  $X^{p,q}$  and have a particularly simple description: the vertices of  $Y^{p,q}$  is the affine line in  $\mathbf{F}_q$ , the finite field of q elements, and quaternions of norm p acting on the affine line as Mobius transformations. These graphs have  $\rho = \max(\lambda_2, -\lambda_n) \leq 2\sqrt{p-1}$ . Each  $Y^{p,q}$  is covered by  $X^{p,q}$ , and each  $X^{p,q}$  is covered by  $X^{p,qr}$  for any integer r. Furthermore the number of vertices of  $X^{p,q}$  tends to infinity for fixed p as q tends to infinity.

#### 6 Fiber Products and Numerical Calculations

The notions of various types of graph products abound in the literature; for many of these products the spectrum is easily written in terms of the spectra of the terms of the product.

One product where this is not the case, and which is not usually mentioned in spectral theory is the fiber product,  $G_1 \otimes_G G_2$  of two graphs  $G_1, G_2$  lying above G (see the appendix for its definition). We are interested in the case that  $G = B_d$ , the directed boquet of d loops; i.e.  $G_i$  are d-regular directed graphs, with a fixed d-coloring of the edges. Then the fiber product, we which abbreviate  $G_1 \otimes_d G_2$ , is a d-regular graph, which is the same as one of the standard graph products where we keep only edge pairs of the same color. We remak that all graphs are graphs over  $B_1$ , and the fiber product  $\otimes_{B_1}$  is one of the stardard products, namely that which takes graphs with adjacency matrices  $A_1, A_2$  and produces one with adjacency matrix  $A_1 \otimes A_2$  (and whose spectrum is therefore the pointwise product of the spectrum of  $A_1$  and  $A_2$ ).

The fiber product  $G_1 \otimes_d G_2$  has certain interesting properties. The first is that it takes two d-regular graphs and produces a larger d-regular graph; it suggests new graphs and stands a chance of being a building block for graphs with, say, small second eigenvalue. Secondly it enables to give a concise description of certain graphs, which is useful in, say, running numerical experiments on them and other "practical" matters. Thirdly when we twist the fibration (i.e. change the coloring) of, say,  $G_2$ , we get different graphs. We give a numerical example of one twisting which has a somewhat surprizingly small second eigenvalue.

The graphs  $X^{p,q}, Y^{p,q}$  are undirected graphs of degree p+1, coming from Cayley graphs with generators  $g_1, g_1^{-1}, \ldots, g_s^{-1}$  with s=(p+1)/2; they should therefore be viewed as graphs over the undirected boquet of loops,  $B_s$  (see the appendix). The Chinese Remainder Theorem and the Hasse principle (i.e. Strong approximation, see [Kne66], or [LPS88] for another proof of surjectivity) easily show that

$$X^{p,q_1q_2} = X^{p,q_1} \otimes_{\mathfrak{s}} X^{p,q_2}$$

for  $q_1, q_2$  relatively prime. This is useful in constructing more simple examples of graphs of this type— aside from  $Y^{p,q}$ , which are quite simple to describe, so are their fiber products

$$Y^{p,q_1} \otimes_s \cdots \otimes_s Y^{p,q_r}$$

for distinct  $q_1, \ldots, q_r$ . That they also have  $\rho \leq 2\sqrt{d-1}$  follows from the fact that they are covered by  $X^{p,q_1\cdots q_r}$ .

Consider graphs of the form  $Y^{p,q_1} \otimes_s Y^{p,q_2}$ . If one changes the fibration  $\pi_2: Y^{p,q_2} \to B_s$  one might wonder what happens to the eigenvalues. In particular, one might expect that having  $\rho \leq 2\sqrt{d-1}$  is a delicate property, and that tampering with the fibration destroys this property. Numerical experiments show that this is not generally the case, and in fact  $\rho$  often decreases when the fibration is changed(!).

Given a cover G of  $B_s$ , with a fixed covering map  $\pi$ , and a cover  $\sigma: B_s \to B_s$ , we write  $G_{\sigma}$  for the cover  $(G, \sigma \circ \pi)$ ; we call  $G_{\sigma}$  a twisting of G. Given two graphs  $G_1, G_2$  over  $B_s$ , a twisted product is any fiber product of twisting of  $G_1, G_2$ . Clearly any such twisted product is isomorphic to a twisted product with  $G_1$  untwisted.

We first consider all the possible twisted products of  $Y^{5,q_1}$  with  $Y^{5,q_2}$ ,  $Y^{5,q_1} \otimes_3 Y^{5,q_2}_{\sigma}$  for appropriate  $\sigma$ . In choosing  $\sigma$  we note that there are 3 orientations to choose and a permutation of  $\{1,2,3\}$  to choose, for a total of 48 possible  $\sigma$ 's. Also the  $Y^{5,q}$  involve

choosing a square root of -1 modulo q, which introduces more possibilities. We claim that there are really only two different graphs that one can generate. We give a table of the largest eigenvalues in absolute value for the two graphs (see table 1).

$q_1 =$	$q_2 =$	$\lambda(Y^{5,q_1}\otimes Y^{5,q_2})$	$\lambda(Y^{5,q_1}\otimes Y^{5,q_2}_\sigma)$	$\binom{5}{q_1}$	$\binom{5}{q_2}$	$\binom{q_1}{q_2}$	n
13	17	-4.3728327	-4.3083943	-1	-1	1	252
13	29	4.3294429	-4.2662239	-1	1	1	420
13	37	4.3854430	4.8793826	-1	-1	-1	532
13	41	-4.3929501	-4.4925255	-1	1	-1	558
13	53	4.4497068	-4.3914762	-1	-1	1	756
13	61	4.4647824	4.4189172	-1	1	1	868
17	29	-4.4442100	4.3346764	-1	1	-1	540
17	37	-4.3804530	4.4278100	-1	-1	-1	684
17	41	4.4536610	4.3783471	-1	1	-1	756
17	53	4.4497068	-4.3914762	-1	-1	1	972

Table 1:  $\lambda$  of twisted products,  $n \le 1000, p = 5 \ (2\sqrt{5} = 4.4721359...)$ .

In this table,  $n = (q_1 + 1)(q_2 + 1)$  is the number of vertices of the graph,  $\lambda$  is the second largest eigenvalue in absolute value, the  $\lambda$ 's are truncated (not rounded) to 7 places (with error  $< 10^{-7}$ ), and  $\sigma$  is any orientation preseving map on  $B_3$  which fixes one edge and transposes the other two. We also warn the reader that these values were generated with a randomized algorithm, so that their correctness is not guarenteed; one can say, however, that these numbers are lower bounds for the second eigenvalue and are very unlikely to be incorrect. Similar remarks about accuracy and correctness hold for all of the tables given here.

In the above table we give the Legendre symbols of the primes involved, in case the reader can spot any patterns. We include some more values, in table 2. This includes a few  $q_1 = 29$  values; it includes  $q_1 = 13$  and  $q_2 \le 157$ ,  $q_2 = 157$  being the first point at which the Legendre symbol  $\binom{q_1}{q_2}$  does not determine whether or not the twist has a smaller second eigenvalue (the author knows of no reason for this bizarre behavior for smaller values of  $q_2$ ); finally we include for  $q_1 = 13,17$  all other  $q_2 < 300$  such that the twist improves the second eigenvalue. The data also suggests that as the number of vertices increases examples of where twisting helps occur less often. We also note that  $Y^{5,173}$  has so poor a second eigenvalue that it persists in its lifting to its products with  $Y^{5,13}, Y^{5,17}$  and their twists.

To explain why there are only two graphs generated, we note the following symmetries. Any automorphism,  $\phi$ , of the standard quaternion algebra mod q,  $\mathbf{H}_q$ , is given as an inner automorphism,  $\phi_a(x) = axa^{-1}$  for some  $a \in \mathbf{H}_q$  (by the Noether-Skolem theorem, see [Pie82, Wei74], or the reader can check all the relevant cases by hand). Recall that  $Y^{p,q}$ 's edges correspond to the set of quaternions, coming in conjugate pairs,  $E_p = \{\alpha \in \mathbf{H} \mid \alpha \overline{\alpha} = p, \quad \alpha \equiv 1 \pmod{2}\}$  (reduced modulo q). Let  $a \in \mathbf{H}_q$  be such that  $\phi_a$  maps  $E_p$  (i.e. its

$q_1 =$	$q_2 =$	$\lambda(Y^{5,q_1}\otimes Y^{5,q_2})$	$\lambda(Y^{5,q_1}\otimes Y^{5,q_2}_\sigma)$	$\binom{5}{q_1}$	$\binom{5}{q_2}$	$\binom{q_1}{q_2}$	n
29	37	4.4530135	4.4328406	1	-1	-1	1140
29	41	-4.4126792	4.5059911	1	1	-1	1260
29	53	-4.4174373	-4.6631110	1	-1	1	1620
13	73	4.4081475	4.4853949	-1	-1	-1	1036
13	89	-4.4391579	-4.4967407	-1	1	-1	1260
13	97	-4.4142135	-4.4493451	-1	-1	-1	1372
13	101	4.4459922	-4.4088410	-1	1	1	1428
13	109	-4.3828926	-4.4246632	-1	1	-1	1540
13	113	4.4610164	4.4536431	-1	-1	1	1596
13	137	4.4511993	-4.4723840	-1	-1	-1	1932
13	149	-4.3981705	-4.4819878	-1	1	-1	2100
13	157	4.4594115	-4.5209753	-1	-1	1	2212
13	173	-4.4533279	-4.4533279	-1	-1	1	2436
13	257	4.4557300	4.4477234	-1	-1	1	3612
17	61	-4.4163100	4.3948404	-1	1	-1	1116
17	109	4.4519048	-4.4417525	-1	1	<sub>3</sub> -1	1980
17	173	-4.4533279	-4.4533279	-1	-1	-1	3132
17	241	-4.4700880	4.4620771	-1	1	-1	4356

Table 2: More  $\lambda$  of twisted products, p = 5 ( $2\sqrt{5} = 4.4721359...$ ).

image reduced modulo q) to itself. Then a determines a twisting,  $Y_a^{p,q}$ , of  $Y^{p,q}$ , via  $\phi_a$ 's twisting of  $E_p$ .

**Proposition 6.1**  $Y_a^{p,q}$  is  $B_s$ -isomorphic to  $Y^{p,q}$ . For any q' we have  $Y^{p,q'} \otimes_s Y_a^{p,q}$  is  $(B_s-)$  isomorphic to  $Y^{p,q'} \otimes_s Y^{p,q}$ .

**Proof** The first isomorphism is given by the map on affine space mod  $q, x \mapsto \tilde{a}x$ , where  $\tilde{a}$  is the Möbius transformation corresponding to a. The second statement follows from the first.

There are 24 automorphisms of  $\mathbf{H}_q$  onto itself which map  $E_p$  to itself, such as  $(i, j, k) \mapsto (j, k, i)$ ,  $\mapsto (-i, -j, k)$ ,  $\mapsto (i, k, -j)$ , etc. Hence of the 48 twists of  $B_3$ , we get only two conceivably non-isomorphic graphs.

In the above we are assuming that we fix a square root,  $\mu_q$  of -1 modulo q, which gives us a bijection  $\nu \colon \mathbf{H}_q \to \mathrm{PGL}(2,q)$  from which we obtain  $Y^{p,q}$ . Using  $-\mu_q$  instead of  $\mu_q$  gives rise to a new bijection,  $\tilde{\nu}$  and new graph  $\tilde{Y}^{p,q}$ . But the new graph is  $B_s$  isomorphic to  $Y_a^{p,q}$  where a comes from the automorphism  $\nu^{-1}\tilde{\nu}$ . So this process does not generate any new graphs.

The above table gives some examples of where the twisted product has smaller eigenvalue than the original product. The author does not have a simple interpretation for the arithmetic meaning of the twisted product, and the fact that they sometimes have  $\rho > 2\sqrt{5}$  (as shown in the table) suggests that they may not have such an interpretation.

We compare these eigenvalues to that of the  $Y^{5,q}$ 's and to random 6 regular graphs. Table 3 gives a list of eigenvalues of  $Y^{5,q}$  for q less than around 1000. To compare these eigenvalues

q =	λ	q =	λ	q =	λ	q =	λ
13	3.5615528	229	-4.3969876	457	4.4213551	757	4.3259582
17	-3.4641016	233	4.4475483	461	-4.4211984	761	4.3729179
29	4.1413361	241	4.3539454	509	-4.3861114	769	4.4575070
37	-4.0000000	257	4.3808292	521	-4.4404421	773	4.4152562
41	-4.1563251	269	-4.4281440	541	-4.4352388	797	4.4625982
53	-4.1563251	277	-4.4506516	557	4.4441941	809	4.3544312
61	-4.1826214	281	-4.4501734	569	-4.3864854	821	-4.4051259
73	4.4081475	293	4.4201881	577	4.4498917	829	-4.4249464
89	4.2659810	313	-4.4431244	593	-4.4108744	853	-4.4502504
97	-4.3141549	337	4.4203043	601	-4.3554867	857	-4.4224954
101	4.2488999	349	4.3650758	613	4.4258625	877	-4.4608373
109	4.2192444	353	4.4570035	617	-4.4024326	881	-4.4300769
113	4.2715228	373	-4.3902568	641	4.4542392	929	-4.4514668
137	4.2254605	389	4.4500992	653	-4.4142135	937	4.4080736
149	4.3415806	397	4.4501133	661	4.4250915	941	-4.4342184
157	4.3247905	401	4.4283013	673	-4.4005898	953	-4.4641787
173	-4.4533279	409	4.4161747	677	-4.4267661	977	-4.3878779
181	4.4207598	421	4.3893717	701	-4.3894433	997	4.4346157
193	4.3142952	433	-4.3823991	709	4.4448742	1009	-4.4239189
197	4.3224143	449	4.4021751	733	-4.4563859	1013	-4.3738233

Table 3:  $\lambda$  of  $Y^{5,q}$   $(2\sqrt{5} = 4.4721359...)$ .

and those of some of their products to random graphs, we sampled 1000 random graphs with a fixed number of vertices, constructed from 3 random permutations, and computed their eigenvalues; here random means picking a random permutation and applying 200 extra random transpositions for good measure, based on the C routine irand48(). This was performed twice for certain small values of vertices (which are = q+1 for small primes  $q \equiv 1 \pmod{4}$ ), to compare random versus the  $Y^{5,q}$ . For the twisted products with  $(q_1, q_2)$  being (13, 17) and (13, 29), we did the experiment once. The results are listed in table 4. These results show that the average spectral radius of a random graph tends to be  $< 2\sqrt{d-1}$ ; at present, it is only known that the average spectral radius is  $\leq 2\sqrt{d-1} + 2\log d + C$  for some constant C (and n sufficiently large depending on d) for most graphs (see [Fri88]), and Alon's conjecture

n	Smallest $\lambda$	$10 h \lambda$	$100 h \lambda$	$550 h \lambda$	10th larg. $\lambda$	Larg. $\lambda$	Avg. $ \lambda $
14	-3.2822140	3.5045810	-3.8625061	-4.3131283	5.2895440	6.0000000	4.3238381
1000	-3.3009647	-3.5196129	-3.8650207	-4.3079389	5.3359077	6.0000000	4.3125449
18	-3.5546727	3.7601589	3.9815828	-4.3569313	5.0594068	6.0000000	4.3550564
-	-3.4615425	3.7424164	-3.9865187	4.3529906	5.0623906	6.0000000	4.3522898
30	-3.7825486	-3.9959083	4.1471119	4.3921495	4.9160898	-5.1557566	4.3984882
	-3.7701586	3.9575758	4.1347205	-4.3796709	4.9616645	6.0000000	4.3995936
38	3.9451714	4.0835724	-4.2027038	-4.4072455	4.8455419	4.9867719	4.4107891
	3.9999286	4.0569757	4.1994161	-4.4008422	-4.7870256	-5.0210640	4.4055494
42	3.9708077	-4.0552115	-4.2282914	4.4110506	-4.8076211	6.0000000	4.4214296
	-3.9781894	4.0704865	-4.2205221	4.4128289	4.8257627	5.0294567	4.4239522
54	-3.9818649	4.1446522	-4.2665110	-4.4177893	4.7635511	5.0450831	4.4259959
	4.0557738	4.1560708	-4.2662609	-4.4189545	4.7904809	4.9765362	4.4328321
252	4.3178301	4.3666513	4.4015240	4.4528272	4.6187277	4.7921238	4.4579061
420	-4.3814378	4.3991200	-4.4227972	4.4577302	4.5681739	4.6886937	4.4622456

Table 4:  $\lambda$  of 1000 random 6-regular graphs  $(2\sqrt{5} = 4.4721359...)$ .

that for fixed d most graphs on n vertices have  $\lambda_2 \leq 2\sqrt{d-1} + \epsilon(n)$  with  $\epsilon(n) \to 0$  as  $n \to \infty$  remains unresolved.

As a rule the fiber product  $Y^{5,q_1}\otimes_s Y^{5,q_2}$  has smaller second eigenvalues than  $Y^{5,q}$  type graphs with comparable number of nodes. Sometimes the twisted fiber products, such as those with  $(q_1,q_2)$  being (13,17),(13,29),(17,41), are much better than comparable  $Y^{5,q}$  graphs. Also, as a rule it seems easy to find 6 regular graphs on q+1 vertices with smaller second eigenvalue than that of  $Y^{5,q}$ , simply by taking enough random graphs, at least when q is small. To beat the fiber products and their twists seems harder, at least in the first two cases. The tables show that at least 10 out of 1000 random graphs sampled beat  $Y^{5,13}\otimes_s Y^{5,17}$ . In looking for graphs with 252 nodes that beat  $Y^{5,13}\otimes_s Y^{5,17}_{\sigma}$ , a search of 100,000 random graphs yielded only four graphs with  $\rho < 4.31$ , three of which beat  $Y^{5,13}\otimes_s Y^{5,17}_{\sigma}$ , which were:

$$-4.2937558, \quad -4.2989024, \quad 4.3081363, \quad -4.3089338 \; .$$

To beat  $Y^{5,13} \otimes_s Y^{5,29}$  and its twist, a search was made among over 1,000,000 random graphs for one with  $\rho < 4.33$ ; none were found. Following a suggestion of Michael Rabin, a heuristic search was made for degree 6 graphs on 420 nodes with  $\rho < 4.33$ : a graph was chosen at random, and each for each of 50,000 iterations 100 random transpositions (in total) were applied to the three permuations from which the graph was constructed; and if the new  $\rho$  decreased, then the transpostions were left, if not the old graph was kept. Such a search yielded one graph with  $\rho < 4.33$ , namely with second eigenvalue

4.3249334,

on the 48,247th iteration (beating  $Y^{5,13} \otimes_s Y^{5,29}$ ). A repeated experiment yielded no graph with  $\rho < 4.33$  for 50,000 iterations. Several variations of this experiment, allowing a "simulated annealing" type step, and with varied number of transpositions per iteration using 50,000 iterations, again yielded no such graphs. Finally, a heuristic search of the above type was tried starting at  $Y^{5,13} \otimes_s Y^{5,29}_{\sigma}$ , introducing one random transposition at a time, to try to improve the second eigenvalue; many graphs with smaller second eigenvalue were found, one as low as

4.2420110,

whose adjacency matrix differed from that of  $Y^{5,13} \otimes_s Y_{\sigma}^{5,29}$  in 147 rows. We also remark that the lower bound in section 3 does not give a very interesting lower bound; namely numerical calculations show that  $T_{6,3}, T_{6,4}$  respectively have valences 3.71899 and 3.94016.

The next smallest possible value of p, p=13, already has many more possibilities for the twisting. For example, the second eigenvalue of  $Y^{13,5} \otimes_7 Y^{13,17}$  was computed as -6.8284271. 33,020 random twistings of the product were made, twisting the permutation of the seven pairs and of the orientations. The following lists the smallest eigenvalue found, the 330th smallest, the 3,302th smallest, the median, the 330th largest, and the largest:

6.5429307, -6.6034627, 6.7329552, 6.9476910, -7.3940794, 8.1223408.

The smallest eigenvalue was attained 33 of the times, and many eigenvalues were repeated around 30 times. This might seem a lot, since  $7!2^7 = 645120$ , but one has to recall all the symmetries present. The above theorem enables us to twist both  $Y^{p,q_1}$  and  $Y^{p,q_2}$  by the 24 automorphism of the quaternion algebra mapping  $E_{13}$  to itself; hence, for any twisting  $\sigma$  on  $q_2$  any twisting  $\phi_a \sigma \phi_b$  gives us the same graph, with  $\phi_a$ ,  $\phi_b$  as described above. When  $\sigma$  is the identity then this yields only 24 twistings, however if  $\sigma$  is very non-arithmetical this could yield conceivably  $24^2$  twistings. Since  $645120/24^2 = 1120$ , it is not surprizing that out of 33,020 random twists many of them were repeated approximately 30 times. The twist which gave the smallest value was the one which took  $\{+1,\ldots,+7\}$  to  $\{-1,-5,-4,7,-2,-6,3\}$ , the  $\pm$  indicating orientation, and the quaternions in order being

$$3+2i, 3+2j, 3+2k, 1+2i+2j+2k, 1+2i+2j-2k, 1+2i-2j+2k, 1+2i-2j-2k$$
.

Nati Linial has suggested to the author that it would be interesting to compare the isoperimetric properties of the two fiber products with different fibrations. It might be a little ambitious, given current technology, to find the best isoperimetric sets, but one could certainly compare the expansion of sets such as balls of a given radius, etc.

#### 7 Remarks

We gather some remarks and suggestions for further research.

We have studied some graphs numerically by calculating their second eigenvalue. In many applications one is more interested in their isoperimetric properties, such as how good expanders they are and how few edges induced subgraphs have. It would be interesting to see if in the above examples small second eigenvalue is correlated to these properties, either provably or by numerical experiments.

The fiber product deserves more study. One thing one could consider is for a fixed d and n (with dn even), consider the family of all d-regular graphs on n nodes over the directed  $B_d$ ,  $\mathcal{G}_{n,d}$  (or perhaps over the oriented undirected  $B_{d/2}$  assuming d even). If one takes the fiber product of all of these graphs, one gets a type of universal graph. Of course, by Galois theory we know such a product will have many connected components and be, for example, a lousy expander. However, perhaps there is some way to choose a subset of such graphs (e.g. a sufficiently fast algorithm, perhaps computing eigenvalues or randomly testing isoperimetric properties), to produce graphs which can be used for certain explicit construction questions. This is an analogue of the Justesen type constructions in coding theory.

From the more abstract point of view, it would be nice to know in what ways covering/Galois theory for graphs can be related to the same for fields. For example, a graph corresponds to a ring and a connected graph corresponds to a field, given the behavior of  $\otimes$  in both theories. Is there a natural ring associated to a graph, so that the ring is a field for connected graphs? Can this be done for certain types of graphs, such as Cayley graphs? Is there an analogue for the notion of modules, exact sequences, etc. in graph theory? Given a finite Cayley graph one can, of course, construct a finite group scheme whose group structure is that of the Cayley graph's underlying group, but then the ring of global functions on the scheme is only a field for the trivial group.

From the point of view of analysis, one of the peculiarities of this theory is due to the fact that the measure  $\mu_1$  is supported at only a finite number of points. One could, of course, work with different measures, such as taking  $\mu = \mu_2 = \mu_1$  to be Lebesgue measure on the edges. One would then get an eigenvalue theory with an infinite number of eigenvalues, spanning all of  $L^2(\mathcal{G}, \mu)$ , etc. One would get an infinite number of eigenpairs, and a similar theory of nodal regions, etc., as well as theorems about the growth rate of the eigenvalues (as in the classical case), holds here. One could study graphs in this way.

In particular, some interest in spectral theory for graphs was due to the possibility of distinguishing non-isomorphic graphs via their spectra (which is not generally possible for the classical graph spectum). If we vary the measures  $\mu_i$  in the Rayleigh quotient, even taking them to have, say, bounded non-zero Radon-Nykodim derivative w.r.t. Lebesgue measure, so that the number of eigenvalues would be infinite, a computer could still, presumably, calculate approximations to some finite number of the eigenvalues. To ask a concrete question, fix an integral constant c > 0; given a graph and a sequence of c edges, consider the measure  $\mu$  which is Lebesgue measure times i+1 on the i-th edge for  $i \leq c$ , and is ordinary Lebesgue measure elsewhere. Given two non-isomorphic graphs, is there a choice of a sequence of c edges in the first so that for any choice of c edges in the second graph the resulting spectra differ somewhere in the first  $n^c$  eigenvalues (by a difference computable in polynomial time)?

Finally, as we suggested in section 4, it would be nice to know more about how the eigenfunctions look, e.g. are their vertices with much larger function values than others, what do their nodal regions look like, etc. In particular, is it true that any d-regular graph with sufficiently many vertices, perhaps satisfying some reasonable properties, has the inradius of at least one of the nodal regions of the second eigenvalue larger than a fixed constant?

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## A Mathematical Background

#### A.1 Shannon's algorithm and VLG's

We recall the Perron-Frobenious theorem. Let A be a non-negative square matrix, whose associated (weighted) graph  $G_A$ , is strongly connected. Then there is a positive eigenvector,  $v_1$ , unique up to scalar; let  $\lambda_1 > 0$  be its eigenvalue. Then all other eigenvalues are  $\leq \lambda_1$  in absolute value. If A is aperiodic, i.e.  $A^n$  is positive for some positive n (i.e. the GCD of all cycles in  $G_A$  is 1), then  $\lambda_1$  is strictly larger than all other eigenvalues in absolute value. We shall call  $\lambda_1(A)$  the (average) valence of A; its base 2 logarithm is known as A's capacity. If v is any non-negative vector with  $Av \leq cv$  for a scalar c, then  $c \geq \lambda_1$ , with strict inequality if(f) Av is strictly less than cv in at least one coordinate; similarly with  $\leq$  and  $\geq$  reversed.

If A is a matrix which is non-negative away from the diagonal,  $G_A$  still being strongly connected, then since A + Ir is non-negative for sufficiently large r, A still has a unique positive eigenvector,  $v_1$ . We again call  $\lambda_1(A)$  the valence, however  $\lambda_1$  need not be postive, and even if it is positive it need not be the largest eigenvalue in absolute value.

If G is the nodal region of some (say Dirichlet) eigenfunction (of a larger graph), then  $A_0(G)$ 's unique positive eigenvector is given by the restriction of the eigenfunction to G. Hence the first Dirichlet eigenvalue of G is the valence of  $A_0(G)$  in the above sense.

For an aperiodic 0-1 matrix, A, the valence and/or capacity measures the number of walks of length k in  $G_A$  for large k. Ditto if A has arbitrary non-negative entries, as long as we weight the walks by the product of the edge weights. A natural generalization of this is to variable length graphs, abbreviated VLG's. A VLG is a graph G = (V, E) along with a positive integral length,  $\ell(e)$ , for each  $e \in E$ . If W(k) denotes the number of walks in G of length  $\leq k$ , length being the sum of the length of the edges, then the valence of G,  $\lambda_1(G)$ , is defined to be the limit of  $W(k)^{1/k}$  as  $k \to \infty$ ; we consider walks of length  $\leq k$  instead of = k to avoid the problems when the GCD of the lengths are > 1, and to facilitate the definition with arbitrary positive edge lengths.

Shannon's algorithm (see [SW49], chaper 1, section 1) for computing the valence (or

capacity) is as follows: let  $Z_G$  be the matrix whose i, j entry is 0 if  $(i, j) \notin E$  and  $z^{\ell}$  where  $\ell$  is the length of (i, j) otherwise, with z a formal parameter. Then  $\lambda_1(G)$  is the reciprocal of the smallest real root of

$$\det(I - Z_G) = 0.$$

More generally, if G has multiple edges and/or positive edge weights associated with the edges, then for each edge of length  $\ell$  and weight w we include a  $wz^{\ell}$  term in  $Z_G$  and the above algorithm yields the valence of G, generalized in the obvious way. It is easy to see that the notion of valence and algorithm generalize for graphs with arbitrary positive edge lengths and with possibly an infinite number of edges (but finite vertex set!).

Aside from arising naturally in information theory (e.g. Shannon's telegraph example), VLG's often give simple ways of working with the capacities of ordinary graphs (see [AFKM86]). Namely, given a graph G = (V, E) and a subset  $V' \subset V$ , the realization of G with vertex set V' is the graph G' = (V', E') where E' consists of the set of walks starting and ending in V' which traverse no V' vertices except at the beginning and end. Each E' edge comes with a weight and length, depending on the weights and lengths of G's edges, in the obvious way. It is not hard to see that G' and G have the same capacity. Typically G' will have an infinite number of edges; this is always the case when the induced subgraph of G on V - V' contains a cycle (assuming the G is strongly connected).

In the case when V' consists of a single vertex, v,  $Z_{G'}$  consists of a single entry, which is a power series,  $P_v^G(z)$ , representing the number of walks from v to itself which don't pass through v in the middle; the  $z^k$  coefficient grows like the valence of the valence of the induced subgraph of G on  $V - \{v\}$ . Shannon's algorithm implies that if  $G_1, G_2$  are two strongly connected graphs,  $v_1, v_2$  vertices of the graphs, and the number of walks from  $v_1$  to itself of a given length is  $\geq$  those of  $v_2$  to itself, then  $\lambda_1(G_1) \geq \lambda_1(G_2)$ ; furthermore strict inequality holds iff the number of walks of some length from  $v_1$  to itself is strictly greater than those of  $v_2$  to itself.

#### A.2 Eigenvalue Comparison Theory

Here we give the proof of theorem 2.3. The proofs are the same as in the classical case, which can be found in [CH53, Gar66].

Proof The max-min principle says that

$$\lambda_k(G) = \max_{\dim(V) = k-1} \min_{f \in V^{\perp}} \mathcal{R}(f),$$

with V taken over all subspaces V of the relevant function space H' of H = H(G) (e.g.  $H_0$  for the Dirichlet eigenvalues, H for the Neumann eigenvalues), and  $\bot$  taken in the  $\mu_1$  inner product. The min-max principle says that

$$\lambda_k(G) = \min_{|F|=f} \max_{f \in F} \mathcal{R}(f),$$

where the max is taken over all subsets F of k mutually perpendicular non-zero vectors. Furthermore, equality holds for f in the above equations if and only if f is an eigenfunction. Both principles are easy to check.

If f is a first eigenfunction (for some connected graph and some boundary conditions), then the edgewise linear function  $\tilde{f}$  whose values on the interior vertices are those of |f| has a Rayleigh quotient  $\leq$  that of f, with strict inequality if there is an edge joining two interior vertices with f values of opposite sign. Furthermore,  $\tilde{f}$  cannot vanish on any interior vertices, for if it does so on a vertex, v, we can assume that it does not vanish on a neighbor of v, and then the function  $f + \epsilon \chi_v$ , where  $\chi_v$  is the characteristic function of v, has smaller Rayleigh quotient for sufficiently small, positive  $\epsilon$ . Hence  $\tilde{f}$  and thus f never vanishes at an interior vertex, and so f's values at all the interior vertices are all positive or all negative.

No two functions with a strict sign on interior vertices can be orthogonal. Thus the first eigenvalue has multiplicity one, and any other eigenfunction has positive and negative values on the interior vertices.

The fact that an eigenfunction restricted to a nodal region is an eigenfunction of that region with appropriate boundary conditions follows by inspecting the corresponding Laplacians. The fact that is the *first* eigenfunction follows from the fact that it has a sign on the interior vertices.

The next statement, that  $\subseteq$  implies  $\le$  for eigenvalues follows by applying the min-max principle to  $G_2$ , with F being the eigenfunctions for  $G_1$  extended by zero. The proof of strict inequality for the first eigenvalue given strict containment follows from the fact that equality in the min-max principle can only hold for f being an eigenfunction, and yet the first eigenfunction is necessarily edgewise linear and has a strict sign on its interior vertices.

Comparing Dirichlet and Neumann eigenfunctions follows from the min-max principle applied to the Neumann problem (i.e. to H as opposed to  $H_0$ ) as in the last paragraph. For strict inequality with the first eigenvalue, note that along any boundary edge a Neumann eigenfunction must be constant. On the other hand a Dirichlet eigenfunction is zero on a boundary vertex and has a strict sign at the connecting interior vertex. So equality cannot hold in the min-max principle unless the first Dirichlet eigenfunction is a Neumann eigenfunction, which can only happen if there is no boundary; on the other hand, when there is no boundary the two clearly agree.

If a k-th eigenfunction, with eigenvalue  $\nu_k$ , has more than k+1 nodal regions, then there is a partition of the graph into k subgraphs,  $G_1, \ldots, G_k$ , such that each  $G_i$  contains a nodal region and  $G_1$  strictly contains a nodal region. If  $f_i$  is the first Dirichlet induced eigenfunction of  $G_i$ , extended by zero outside  $G_i$ , then there exists a linear combination  $f = a_1 f_1 + \cdots + a_k f_k$  of the  $f_i$  such that  $a_1 \neq 0$  and f is orthogonal to the first k-1 eigenfunctions. But the Rayleigh quotients of each  $f_i$  are  $\leq \nu_k$ , and that of  $f_1$  is  $< \nu_k$  by the above; an easy calculation then shows that the Rayleigh quotient of f is  $< \nu_k$ , which is impossible.

The statement about changing  $|a_w-1|$ , the edge lengths, or adding edges, follows from the fact that increasing them does not decrease the Rayleigh quotient evaluated at any function.

The statement about Dirichlet induced eigenvalues of disjointly contained  $G_i$ 's in G follows from the min-max principle. The statement about Neumann induced follows from the fact that if in the integral  $\int |\nabla f|^2 d\mu_2$  we allow f to be discontinuous over a discrete set of points, S, and do not integrate through these points, then the Rayleigh quotient does not increase and the class of functions considered is not decreased; when S is taken to be the boundary

#### A.3 Galois Theory

We recall some notions from covering theory and Galois theory for graphs. The connection of graph covers (or "factors") with the spectum appears frequently in the literature (see, for example, [CDS79]). The notion of Galois theory appears less often in graph theory, but is well known to people studying analysis on trees (e.g. p-adic Lie groups), combinatorial group theory, etc., and probably appears in the literature in various places, at least implicitly, in say, [MKS66, CDS79, Ser80, GT87, DD89, Joh90]. It is a very simple Galois theory, and so we describe it here. It also gives rise to the notion of a fiber product<sup>3</sup>, which can be used to generate some simple and interesting new graphs.

In what follows a graph, G, is the geometric realization of a graph; so G is a metric space, having a distinguished set of points, V, set of subsets, E, etc. Working this way one can give concise definitions that work for all graphs, including those with multiple edges and self-loops. A directed graph is a graph where each edge, being isomorphic to a real interval, is given an orientation. A morphism of graphs is a metric preserving map; for directed graphs it should also preserve the orientations of the edges. The definitions also work for graphs with boundaries.

A morphism,  $\phi: G_1 \to G_2$ , is a covering map if it is locally an isomorphism. A cover of G is a collection  $(H, \phi)$  of a graph and a covering map; if G is connected then the number of inverse images of any point is the same, and is called the *degree* of the cover, denoted [H:G]. Covering induces a partial order on the set of graphs, namely  $H \geq G$  if there exists a covering map,  $\phi: H \to G$ .

For any graph G there is a universal cover or a "largest cover,"  $(\widetilde{G}, \phi)$ .  $\widetilde{G}$  is a tree, and its vertices are just the set of "non-backtracking" walks for a fixed vertex,  $v \in G$ .

For graphs  $G_1, G_2$  lying above a graph G, i.e. which are given fixed homomorphisms  $\phi_i: G_i \to G$ , the *fiber product*,  $G_1 \otimes_G G_2$  is given as the set of pairs,  $(x_1, x_2), x_i \in G_i$ , such that  $\phi_1(x_1) = \phi_2(x_2)$ , viewed as a graph in the obvious way. If  $G_1$  and  $G_2$  are covers of G, then  $G_1 \otimes_G G_2$  is a cover of  $G_1$  and of  $G_2$ .

The above notion coincides with the abstract notion of fiber product. Recall that for a fixed graph G one can consider the set of all graphs lying over G, and for any two such graphs,  $G_1, G_2$ , with fixed homomorphisms  $\phi_i \colon G_i \to G$ , one considers the set of "G (equivariant) morphisms," i.e. morphisms  $f \colon G_1 \to G_2$  such that  $\pi_2 \circ f = \pi_1$ . This is the category  $\mathbf{Gr}_G$ . The fiber product is by definition an H lying over  $G_1, G_2, G$  which is universal in that every other H' lying over the three has its projections factoring in a unique way through H; clearly it exists and is the above product.

<sup>&</sup>lt;sup>3</sup>Perhaps "coupled product" would be more suggestive to, say, people working in Markov chains. Indeed, the walk on the fiber product,  $G_1 \otimes_G G_2$ , viewed as a Markov chain, is just the simultaneous walk on the chains  $G_1$  and  $G_2$ , coupled according to G.

We leave it as an exercise to the reader to give these definition in terms of traditional graph theory. The reader should remark that in figure 1, with a graph G having a self-loop,  $G_1$  is a cover of G and  $G_2$  isn't. We remark that for directed graphs, the fiber product

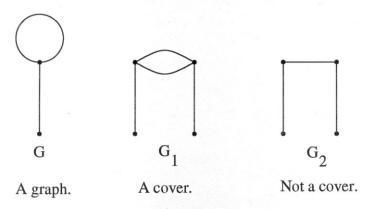


Figure 1: Graph covers for undirected graphs with self-loops.

with G the trivial graph (consisting of a vertex and one self-loop) is just one of the standard products.

To speak of a Galois theory, we want to view the vertices of the universal cover, T, of G as a Cayley graph. This predicates the graph being d-regular for some d; we start by describing the situation for directed graphs.

Let  $B_d$  be the boquet of d loops, i.e. the directed graph with one vertex and d edges (self-loops),  $E_d = \{e_1, \dots e_d\}$ . A cover, G, of  $B_d$  is just a d-regualr graph with a fixed d-coloring of the edges (each vertex having d differently colored edges starting and ending at it). Of course, by Hall's theorem any d-regular digraph G has a covering map to  $B_d$ . The vertices of the universal cover,  $T_d$ , of  $B_d$ , can be identified with the free group on the d elements,  $E_d$ , once we fix a vertex r of  $T_d$ , and  $T_d$  is the Cayley graph of this free group with generators  $E_d$ . For any cover, G, of  $B_d$ , and any vertex v of G, there is a unique covering map,  $\phi: T_d \to V$  with  $\phi(r) = v$ . The set  $H_v = \phi^{-1}(v)$  is a subgroup of  $T_d$ , and G is the quotient graph  $T_d/H_v$  (i.e. graph of left cosets). For any other vertex, v', of G,  $H_v$  and  $H_{v'}$  are conjugate subgroups.

The graph G is called Galois or normal if  $H_v$  is a normal subgroup, equivalently if  $H_v$  does not depend of v, or equivalently G is the Cayley graph of some group. All the stardard results of Galois type theories hold here. For example, any finite graph, G, has a finite Galois cover, i.e. is the quotient of a finite Cayley group with  $\leq n^n$  vertices, where n is the number of vertices of G; indeed,  $H_v$  contains the normal subgroup

$$\bigcap_{[g]\in T/H_v}gH_vg^{-1}$$

which is clearly of index at most  $n^{n-1}$  in  $H_v$ . The reader can compare this Galois theory to that of first year algebra— a graph translates to a ring, a connected graph to a field, covers to field extensions,  $\otimes$  translates to  $\otimes$ , etc. One can check that some well known identities in

field theory, such as

$$L \otimes_K L = \bigoplus_{\sigma \in \operatorname{Gal}(L/K)} L^{\sigma}$$

for Galois extensions L/K, hold for graphs as well (in fact, this identity points out why graphs correspond to rings,  $\otimes$  to  $\otimes$ , etc.).

For example, if H is a Galois cover of a graph G, then Gal(H/G) is just the set of G-equivariant automorphisms from H to itself. If we fix an  $h_0 \in H$  and consider its conjugates  $h_1, \ldots, h_{d-1}$ , i.e. the other elements of  $\pi^{-1}(\pi(h))$ , then each each  $h_i$  gives rise to a unique  $\sigma \in Gal(H/G)$  with  $\sigma(h_0) = h_i$ .

If  $G_1, G_2$  are connected Galois covers of a connected graph, G, then any connected component, H, of  $G_1 \otimes_G G_2$  is a graph covering the former graphs; it is minimal in the sense that any graph which is a G-equivariant cover of  $G_1$  and of  $G_2$  is also a cover of H. Thus a "max" operation exists for the covering order in the subcategory  $\mathbf{Gal}_G$  of Galois covers of  $\mathbf{Gr}_G$ .

If A is a group acting freely on a graph, G, then G/A is a graph, and is covered by G. The point here, of course, is that "freely," by definition, means without fixed points on the geometric realization of G for non-trivial  $a \in A$ , which the reader can check is enough to guarentee a nice quotient (i.e. "regular quotient," as in, say, [GT87]).

We remark that not all aspects of field theory have graph counterparts in Galois theory. For example, we don't know in what naturally arising sense a graph cover can be allowed to be "ramified." Also some of the above has analogues for graphs which are not d-regular (but we cannot hope to realize the universal cover as a Cayley graph in any nice sense).

In spectral theory one is almost always interested in undirected graphs. For undirected graphs there are two typical Galois cases arising. One is Cayley graphs with generators  $\{g_1, \ldots, g_d\}$  which are involutions, i.e.  $g_i^2 = 1$ . The second is Cayley graphs with generators  $\{g_1, g_1^{-1}, \ldots, g_s^{-1}\}$ , d = 2s. Both can be treated as directed graphs (i.e. by thinking of each undirected edge as two directed edges), restricting the covers to  $B_d$  to graphs in which the identity relations implied hold at each vertex of the cover. For example, in the former case this means that following two edges which map to the same  $B_d$  edge leads back to the same vertex. In the latter situation  $B_d$ 's edges come in pairs, and following an edge and then its paired edge leads back to the same vertex; alternatively, we can view these graphs as graphs over over the undirected boquet of s = d/2 loops,  $B_s$ , where each edge comes with an orientation; while an edge can be traversed in either direction, the orientation is a bookkeepping device which tells us how to take fibered products. (We rigorize this by defining the category of undirected oriented graphs, requiring morphism to preserve orientation, etc.) We shall use the notion of oriented undirected graphs, because it simplifies the discussion and notation; covers of  $B_s$  are always 2s-regular undirected graphs.

One should expect, as in all Galois/covering theories, to be able to study some aspects of a graph by considering its covers (e.g. section 4). For studying the spectrum, if H is a cover of G, of finite degree, then G's eigenpairs pull back to eigenpairs of H; H's eigenvectors push forward to vectors on G by summing over the inverse images of points in G, and eigenvectors of H not arising from pullbacks of eigenvectors on G are precisely those which push forward

to 0. Covers of infinite degree present the unpleasantness that eigenvectors of the base graph do not lift to eigenvectors of finite  $L^2$  norm on the cover; relations between the spectra are not as exact as with finite covers. Thus it seems better in some cases to work with finite covers; the universal cover does not exist in this category, but, as usual, should be regarded as the inverse limit of all finite covers.

The same theory yields a covering/Galois theory for t-uniform hypergraphs (i.e. each  $e \in E$  is a set of vertices of fixed size, t) as quotients of hypertrees (see [Fri91]). This situation is a good example of where spectral aspects of the universal cover, the hypertree, which are relatively easy to analyze, are not known to have much bearing on questions about explicitly constructing their finite counterparts (in a way which yields small eigenvalues, good isoperimetric properties, etc.).

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